

SET 11

10.1)

Given;

$$\nabla^2 V + \frac{\partial}{\partial t}(\vec{\nabla} \cdot \vec{A}) = -\rho/\epsilon$$

$$\nabla^2 \vec{A} - \mu\epsilon \frac{\partial^2 \vec{A}}{\partial t^2} - \vec{\nabla}(\vec{\nabla} \cdot \vec{A} + \mu\epsilon \frac{\partial V}{\partial t}) = -\mu\vec{J}$$

The equation for the vector potential is automatically satisfied if we choose;

$$L = \vec{\nabla} \cdot \vec{A} + \mu\epsilon \frac{\partial V}{\partial t}$$

Substitute for  $\vec{\nabla} \cdot \vec{A} = L - \mu\epsilon \frac{\partial V}{\partial t}$  into the equation for the scalar potential to obtain;

$$\nabla^2 V - \mu\epsilon \frac{\partial^2 V}{\partial t^2} + \frac{\partial L}{\partial t} = -\rho/\epsilon$$

10.4)

Given;

$$V = 0 \text{ and } \vec{A} = A_0 \sin(kx - \omega t) \hat{y}$$

Then find the fields;

$$\vec{E} = -\vec{\nabla} V - \frac{\partial \vec{A}}{\partial t} = \omega A_0 \cos(kx - \omega t) \hat{y}$$

$$\vec{B} = \vec{\nabla} \times \vec{A} = k A_0 \cos(kx - \omega t) \hat{z}$$

Substitute into Maxwell's equations to find that;

$$\vec{\nabla} \cdot \vec{E} = 0$$

$$\vec{\nabla} \cdot \vec{B} = 0$$

$$\vec{\nabla} \times \vec{E} = -\frac{\partial \vec{B}}{\partial t}$$

$$\vec{\nabla} \times \vec{B} = \mu\epsilon \frac{\partial \vec{E}}{\partial t}$$

To satisfy the equations we need the dispersion relation;

$$k = \omega \sqrt{\mu\epsilon}$$

10.7)

For the Lorentz gauge we need

$$\vec{\nabla} \cdot \vec{A} + \mu\epsilon \frac{\partial V}{\partial t} = 0$$

Choose a gauge transformation;

$$\vec{A} \rightarrow \vec{A} + \vec{\nabla}\Lambda$$

$$V \rightarrow V - \frac{\partial \Lambda}{\partial t}$$

Thus obtain the Lorentz gauge;

$$\vec{\nabla} \vec{A} + \mu\epsilon \frac{\partial V}{\partial t} = \nabla^2 \Lambda - \mu\epsilon \frac{\partial^2 \Lambda}{\partial t^2} = 0$$

We find a solution to a specific gauge which makes the above equation vanish. We could choose  $V = 0$  and obtain an electric field from  $\frac{\partial \Lambda}{\partial t}$ . We cannot choose  $\vec{A} = 0$  because  $\vec{\nabla} \times \vec{\nabla}\Lambda = 0$  which makes the magnetic field always vanish.

10.10)

The vector potential has the form;

$$\vec{A} = \frac{\mu}{4\pi} \int d\tau' \frac{\vec{J}(\vec{r}', t')}{|\vec{r}' - \vec{r}|}$$

In the above we use the expression for the retarded time,  $t' = t - (|\vec{r}' - \vec{r}|)/c$ . We also note that  $\int d\vec{a} \cdot \vec{J} = I = kt' = kt - kr/c$ , and that  $d\tau' = d\vec{a} \cdot d\vec{l}$  where  $\vec{l}$  points in the direction of the current. The integral around the current loop can be divided into 4 components and each integral evaluated.

$$\begin{aligned} & \int_0^\pi I/r b d\phi \\ & - \int_0^\pi I/r a d\phi \\ & \int_{-a}^{-b} I/r b dx \\ & \int_b^a I/r b dx \end{aligned}$$

Substitute for  $I$  as above and complete the integrals. Collecting terms we find at the origin;

$$\vec{A} = \frac{\mu kt}{2\pi} \ln(b/a) \hat{x}$$

In this case  $V = 0$  so that;

$$\vec{E} = -\frac{\partial \vec{A}}{\partial t} = -\frac{\mu k}{2\pi} \ln(b/a) \hat{x}$$

However  $\vec{B}$  cannot be obtained unless we know more about the spatial dependence of the vector potential.